# **Eigenproblems in Resonant MEMS Design**

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#### What are MEMS?



#### **RF MEMS**



Microguitars from Cornell University (1997 and 2003)

- MHz-GHz mechanical resonators
- Uses:
  - RF signal processing (better cell phones)
  - Sensing elements (e.g. chemical sensors)
  - Really high-pitch guitars

## **Micromechanical filters**



Your cell phone is already mechanical!

- Uses a quartz surface-acoustic wave (SAW) filter
- Can do better using MEMS
  - MEMS filters can be placed on-chip
  - Versus SAWs: smaller, lower power

Success  $\implies$  "Calling Dick Tracy!"

# Damping

- Want to minimize damping
  - Measure by "quality of resonance"

$$Q = \frac{|\omega|}{\mathrm{Im}(\omega)}$$

- Electronic filters have too much
- Understanding of damping in MEMS is lacking
- Several sources of damping
  - Anchor loss
  - Thermoelastic damping
  - Fluid damping
  - Material losses

# Damping

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#### **Example: Disk anchor loss**



- SiGe disk resonators built by E. Quévy
- Axisymmetric model with bicubic mesh, about 10K nodal points

# **Perfectly matched layers**

- Model half-space with a perfectly matched layer
  - Complex coordinate change  $x \mapsto z(x; \omega)$
  - Apply a complex coordinate transformation
  - Generates a non-physical absorbing layer
- Idea works with general linear wave equations
  - First applied to Maxwell's equations (Berengér 95)
  - Similar idea introduced earlier in quantum mechanics (*exterior complex scaling*, Simon 79)















Clamp solution at transformed end to isolate outgoing wave.

### **Choice of transformations**

- $\checkmark$  Generally z depends nontrivially on  $\omega$ 
  - Needed for frequency-independent attenuation
  - Common choice is

$$\frac{dz}{dx} = 1 - \sigma(x)/k$$

- What if we use a fixed transformation?
  - Can choose to absorb well over finite  $\omega$  range
  - Solve a *linear* eigenvalue problem
  - Amounts to rational approx of true radiation condition (in discrete case)

#### **Behavior with fixed transformations**



• Start with  $(K - \omega^2 M)u = e_1$ 

#### **Behavior with fixed transformations**



Schur complement to eliminate PML unknowns

#### **Behavior with fixed transformations**



Compare last coefficient with exact (discrete) BC

# **Complex symmetry**

Finite element equations (forced vibration) are

$$-\omega^2 M u + K u = F$$

where M and K are complex symmetric.

- Row and column eigenvectors are transposes
- Second-order accuracy with modified Rayleigh quotient:

$$\theta(v) = (v^T K v) / (v^T M v)$$

- Can have  $v^T M v \approx 0$ 
  - Propagating modes (continuous spectrum)
  - Not the modes of interest for resonators

### Q variation



- Small geometry variation  $\implies$  large damping variation
- Solid line is simulated; dots are measured

# **Effect of varying film thickness**



 Sudden dip in Q comes from an interaction between a (mostly) bending mode and a (mostly) radial mode

### **Model reduction**

Would like a reduced model which

- Preserves second-order accuracy for converged eigs
- Keeps at least Arnoldi's accuracy otherwise
- Is physically meaningful

Idea:

- Build an Arnoldi basis V
- Double the size:  $W = \operatorname{orth}([\operatorname{Re}(V), \operatorname{Im}(V)])$
- Use W as a projection basis
- Resulting system is still a Galerkin approximation with real shape functions for the continuum PML equations

#### **Example: Disk resonator response**



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# **Thermoelastic damping (TED)**

u is displacement,  $T = T_0 + \theta$  is temperature

$$\sigma = C\epsilon$$

$$\rho u_{tt} = \nabla \cdot \sigma$$

$$\rho c_v \theta_t = \nabla^2 \theta$$

Second-order mechanical + first-order thermal equation

# **Thermoelastic damping (TED)**

u is displacement,  $T = T_0 + \theta$  is temperature

 $\sigma = C\epsilon - \beta\theta 1$   $\rho u_{tt} = \nabla \cdot \sigma$  $\rho c_v \theta_t = \nabla^2 \theta - \beta T_0 \operatorname{tr}(\epsilon_t)$ 

- Second-order mechanical + first-order thermal equation
- Temperature change causes stress (thermal expansion)
- Volumetric strain rate causes thermal fluctuations

# **Thermoelastic damping (TED)**

Non-dimensionalized equation:

$$\sigma = \hat{C}\epsilon - \xi\theta 1$$
  

$$u_{tt} = \nabla \cdot \sigma$$
  

$$\theta_t = \eta \nabla^2 \theta - \operatorname{tr}(\epsilon_t)$$

- **•** Typical MEMS scales:  $\xi$  and  $\eta$  small
- **Perturbation about**  $\xi = 0$  is effective

#### **Perturbation computation**

Discrete time-harmonic equations:

$$-\omega^2 M_{uu}u + K_{uu}u + K_{ut}\theta = 0$$
$$i\omega D_{tt}\theta + K_{tt}\theta + i\omega D_{tu}u = 0$$

Approximate  $\omega$  by perturbation about  $K_{u\theta} = 0$ :

$$-\omega_0^2 M_{uu} u_0 + K_{uu} u_0 = 0$$
$$i\omega_0 D_{\theta\theta} \theta_0 + K_{\theta\theta} \theta_0 + i\omega_0 D_{tu} u_0 = 0$$

Choose  $v: v^T u_0 \neq 0$  and compute

$$\begin{bmatrix} (-\omega_0^2 M_{uu} + K_{uu}) & -2\omega_0 M_{uu} u_0 \\ v^T & 0 \end{bmatrix} \begin{bmatrix} \delta u \\ \delta \omega \end{bmatrix} = \begin{bmatrix} -K_{u\theta} \theta_0 \\ 0 \end{bmatrix}$$

# **Comparison to Zener's model**



- Good match to Zener's approximation for TED in beams
- Real and imaginary parts after first-order correction agree to about three digits with Arnoldi

# Conclusions

- MEMS resonator simulations give interesting problems
- Damped resonators  $\implies$  nonlinear eigenproblems
  - Introduce auxiliary variables to get exact or approximate linear problem
  - There's still useful structure in non-Hermitian problems!
- References:
  - Bindel and Govindjee. "Elastic PMLs for Resonator Anchor Loss Simulation." (IJNME, to appear)
  - HiQLab home page:

www.cs.berkeley.edu/d̃bindel/hiqlab/