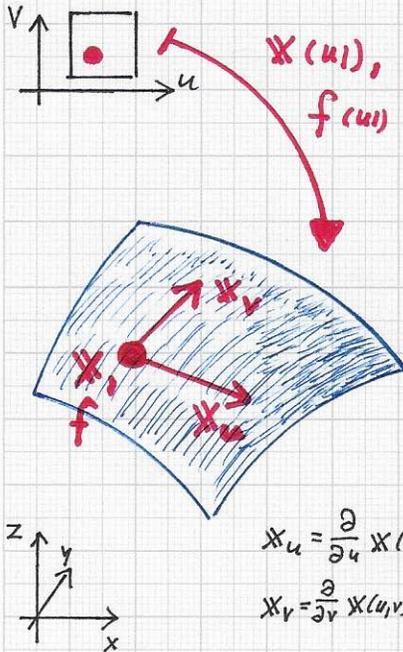


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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

Laplacian eigenfunctions:



Prior to using the Laplacian operator  $\Delta$  to compute Laplacian eigenvalues in the discrete, functional, trivariate setting involving a simple unit voxel grid, it is important to note the differences between our driving problem (computing Laplacian eigenvalue features for unit voxel grids with associated max values) and Laplace-Beltrami eigenvalues that

arise when dealing with a function defined on a curved, 2D manifold domain (or curved, 3D manifold domain) embedded in 3D (4D) space. We sketch the Laplace-Beltrami operator for a function  $f$  on a curved parametric surface  $X(u)$ .

Parametric surface setting for generalized Laplace-Beltrami operator. The function  $f$  to be analyzed has a curved surface as its domain; thus both the surface  $X$  and  $f$  are functions of  $u$  and  $v$ , i.e.,  $X = X(u)$  and  $f = f(u)$ .

Beltrami defines the SECOND differential parameter of  $f = f(u)$  in terms of the curvilinear coordinates  $u$  and  $v$  of a parametric surface  $X = X(u)$ :

$$\Delta f = \frac{1}{L} \left( (Gf_u - Ff_v) / L \right)_u + \left( (Ef_v - Ff_u) / L \right)_v \Big/ L.$$

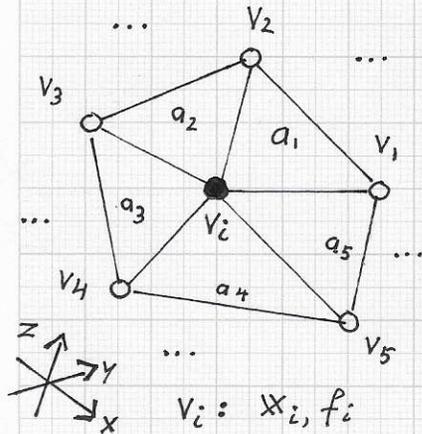
Since  $f$  is a "function on a curved surface," the local differential geometry of the surface at a point  $X$  must be determined and used to compute the generalized  $\Delta f$  value at  $X$ . As the  $\Delta$  operator is generally defined as  $\Delta f = \text{div grad } f$ , one must use the partial derivatives  $X_u$  and  $X_v$  at  $X$  and use the partial derivative vectors to define the "grad" field subjected to the "div" operator.

The "subscripts"  $u$  and  $v$  refer to differentiation with respect to  $u$  and  $v$ , respectively. Further,  $E = \langle X_u, X_u \rangle$ ,  $F = \langle X_u, X_v \rangle$ ,  $G = \langle X_v, X_v \rangle$ , and  $L = (EG - F^2)^{1/2}$ . (Dirk Struik, Lectures on Classical Differential Geometry)

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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

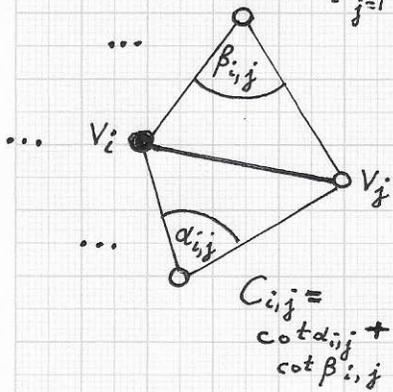
Laplacian eigenfunctions: Beltrami's definition of the second differential parameter of a function  $f = f(u)$



$v_i: x_i, f_i$

$v_j: x_j, f_j, j=1...5$

$a_j: \text{triangle area}$   
 $\Rightarrow A_i = \sum_{j=1}^5 a_j$



$C_{i,j} = \cot \alpha_{i,j} + \cot \beta_{i,j}$

on a curved surface, leading to the general Laplace-Beltrami operator, can be used to analyze a function  $f$  when it is known only at points  $x_i$  in 3D space with associated function values  $f_i$ ; commonly, a simple triangle mesh is used to define a continuous, piecewise linear, domain for the function to be analyzed. Given a set  $\{x_i, f_i\}$  and a mesh, one can locally define many meaningful estimates of  $\Delta f(x_i)$ . In this case, a parametric map  $u \mapsto x(u), u \mapsto f(u)$  is not known. Triangle meshes are commonly used in many applications, and various approximation schemes exist to estimate  $\Delta f$  at mesh vertices  $v_i$  when  $f$ -values are known only at these vertices.

One standard approximation of  $\Delta f(v_i)$  is:

$$\Delta(v_i) = \frac{3}{2A_i} \sum_{j=1}^5 C_{i,j} (f_i - f_j).$$

Local data needed to estimate  $\Delta f$  at a triangle mesh vertex using a commonly used scheme that considers geometrical information (areas and angles) and functional data ( $f$ -values at vertices).

This approximation formula can be derived by viewing  $f$  as a piecewise-defined continuous

Formulae must be able to handle mesh boundaries.

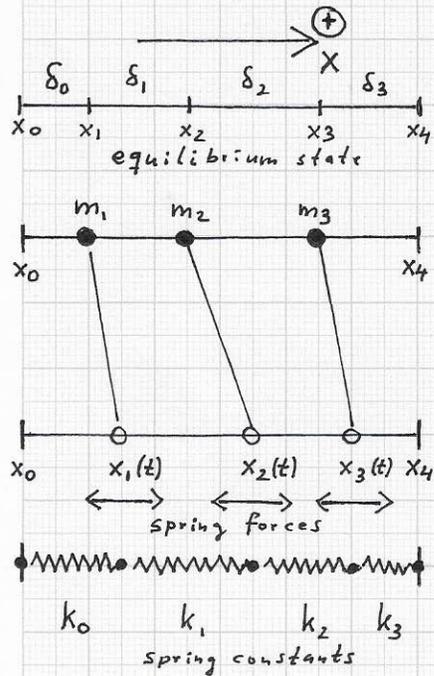
function on the triangle mesh. (Keenan Crane et al., Digital Geometry Processing with Discrete Exterior Calculus, course notes, ACM SIGGRAPH 2013)

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■ OBJECT AND MATERIAL EIGENFUNCTIONS - cont'd.

• Laplacian eigenfunctions:

We are merely interested in the discrete, functional case when considering  $\Delta f$ . More specifically, we only have to define and compute  $\Delta f$ -values for centers (or corners) of voxels with associated masses defining a segment/object/material in a 3D scan. Once we have  $\Delta f$ -values for all voxel centers (or corners), we can set up the eigenvalue problem  $\Delta f = -\lambda f$ . It is helpful to consider the 1D case of of a harmonic oscillator with point masses connected by springs as a linear chain (see illustration, left). This is the set-up:



1D example of a mass-spring model used to derive the eigenvalue problem that defines the oscillating systems eigenmodes/frequencies.

- 3 point masses  $m_1, m_2, m_3$  are given.
- The 3 points' positions in rest state are  $x_1, x_2, x_3$ .
- The "wall end points" are  $x_0$  and  $x_4$ .
- The rest state spacing is given by 
$$\delta_i = x_{i+1} - x_i, \quad i = 0 \dots 3.$$
- Point pairs  $x_i$  and  $x_{i+1}$  are connected by a spring with spring constants  $k_i$ .
- The time-dependent positions of the three "free mass points" are  $x_i(t)$ ,  $i = 1 \dots 3$ .
- Newton's and Hooke's laws describe the systems behavior via second-order differential equations:  $m \ddot{x}(t) = -kx(t)$ .

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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

Laplacian eigenfunctions: Considering the illustration and parameter explanations from the previous page, the system must satisfy these equations:

Derivation of  
Linear system of  
second-order  
differential  
equations  
describing the  
mass-spring  
system of point  
masses on a  
"Linear chain"

$$\begin{aligned} m_1 \ddot{x}_1(t) &= -k_0 (x_1(t) - x_0(t)) - \delta_0 \\ &\quad + k_1 (x_2(t) - x_1(t)) - \delta_1 \\ &= k_0 x_0(t) + (-k_0 - k_1) x_1(t) + k_1 x_2(t) + k_0 \delta_0 - k_1 \delta_1 \end{aligned}$$

$$\begin{aligned} m_2 \ddot{x}_2(t) &= -k_1 (x_2(t) - x_1(t)) - \delta_1 \\ &\quad + k_2 (x_3(t) - x_2(t)) - \delta_2 \\ &= k_1 x_1(t) + (-k_1 - k_2) x_2(t) + k_2 x_3(t) + k_1 \delta_1 - k_2 \delta_2 \end{aligned}$$

$$\begin{aligned} m_3 \ddot{x}_3(t) &= -k_2 (x_3(t) - x_2(t)) - \delta_2 \\ &\quad + k_3 (x_4(t) - x_3(t)) - \delta_3 \\ &= k_2 x_2(t) + (-k_2 - k_3) x_3(t) + k_3 x_4(t) + k_2 \delta_2 - k_3 \delta_3 \end{aligned}$$

⇒

$$\begin{aligned} m_i \ddot{x}_i(t) &= k_{i-1} x_{i-1}(t) + (-k_{i-1} - k_i) x_i(t) \\ &\quad + k_i x_{i+1}(t) \\ &\quad + k_{i-1} \delta_{i-1} - k_i \delta_i, \quad i=1 \dots 3, \end{aligned}$$

where  $x_0(t) = x_0$ ,  $x_4(t) = x_4$ .

The resulting inhomogeneous system is solved with initial conditions  $\mathbb{x}(t_0) = \mathbb{x}_0$  and  $\dot{\mathbb{x}}(t_0) = \dot{\mathbb{x}}_0$ .

In matrix notation, this linear system is:

$$\begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix} \begin{bmatrix} \ddot{x}_1(t) \\ \ddot{x}_2(t) \\ \ddot{x}_3(t) \end{bmatrix} = \begin{bmatrix} -k_0 - k_1 & k_1 & 0 \\ k_1 & -k_1 - k_2 & k_2 \\ 0 & k_2 & -k_2 - k_3 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix} + \begin{bmatrix} k_0 x_0 + k_0 \delta_0 - k_1 \delta_1 \\ k_1 \delta_1 - k_2 \delta_2 \\ k_2 x_4 + k_2 \delta_2 - k_3 \delta_3 \end{bmatrix}$$

$$\underline{M \ddot{\mathbb{x}}(t) = K \mathbb{x}(t) + \mathbb{lb}}$$

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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

Laplacian eigenfunctions: In our application - feature value computation for volumetric object segments defined by voxels with voxel-specific masses - we can set all spring constant values  $k_i$  and all spacing values  $\delta_i$  to 1. In this setting, the inhomogeneous system is

$$\begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix} \ddot{x}(t) = \begin{bmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & 1 & -2 \end{bmatrix} x(t) + \begin{bmatrix} x_0 \\ 0 \\ x_4 \end{bmatrix}$$

Eigenvectors:

$$e_1 = (e_{1,1}^x, e_{1,1}^y, e_{1,1}^z)^T$$

$$e_2 = (e_{2,1}^x, e_{2,2}^y, e_{2,2}^z)^T$$

$$e_3 = (e_{3,1}^x, e_{3,2}^y, e_{3,3}^z)^T$$

Eigenvalues:

$$\lambda_1 = \dots, \lambda_2 = \dots, \lambda_3 = \dots$$

$$M \ddot{x}(t) = K x(t) + lb$$

(mass)                      (stiffness)

The actual eigenvalue problem we are interested in concerns the homogeneous system:

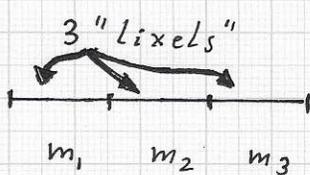
$$\ddot{x}(t) = M^{-1}K x(t)$$

$$\Leftrightarrow \ddot{x}(t) - M^{-1}K x(t) = 0$$

Considering our example with 3 mass points,

the matrix we must analyze is  $M^{-1}K$ :

$$M^{-1}K = \begin{pmatrix} 1/m_1 & 0 & 0 \\ 0 & 1/m_2 & 0 \\ 0 & 0 & 1/m_3 \end{pmatrix} \begin{pmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & 1 & -2 \end{pmatrix} = \begin{pmatrix} -2/m_1 & 1/m_1 & 0 \\ 1/m_2 & -2/m_2 & 1/m_2 \\ 0 & 1/m_3 & -2/m_3 \end{pmatrix}$$



$\Rightarrow$  3 eigenvalues for 3-lixel segment

In this simple "1D segment scenario", one associates masses  $m_1$ ,  $m_2$  and  $m_3$  with line elements ("lixels") of length 1, and the 3 eigenvalues define a feature for this segment. (Indexing of and connectivity between mass points  $\Rightarrow K$  is tridiagonal.)