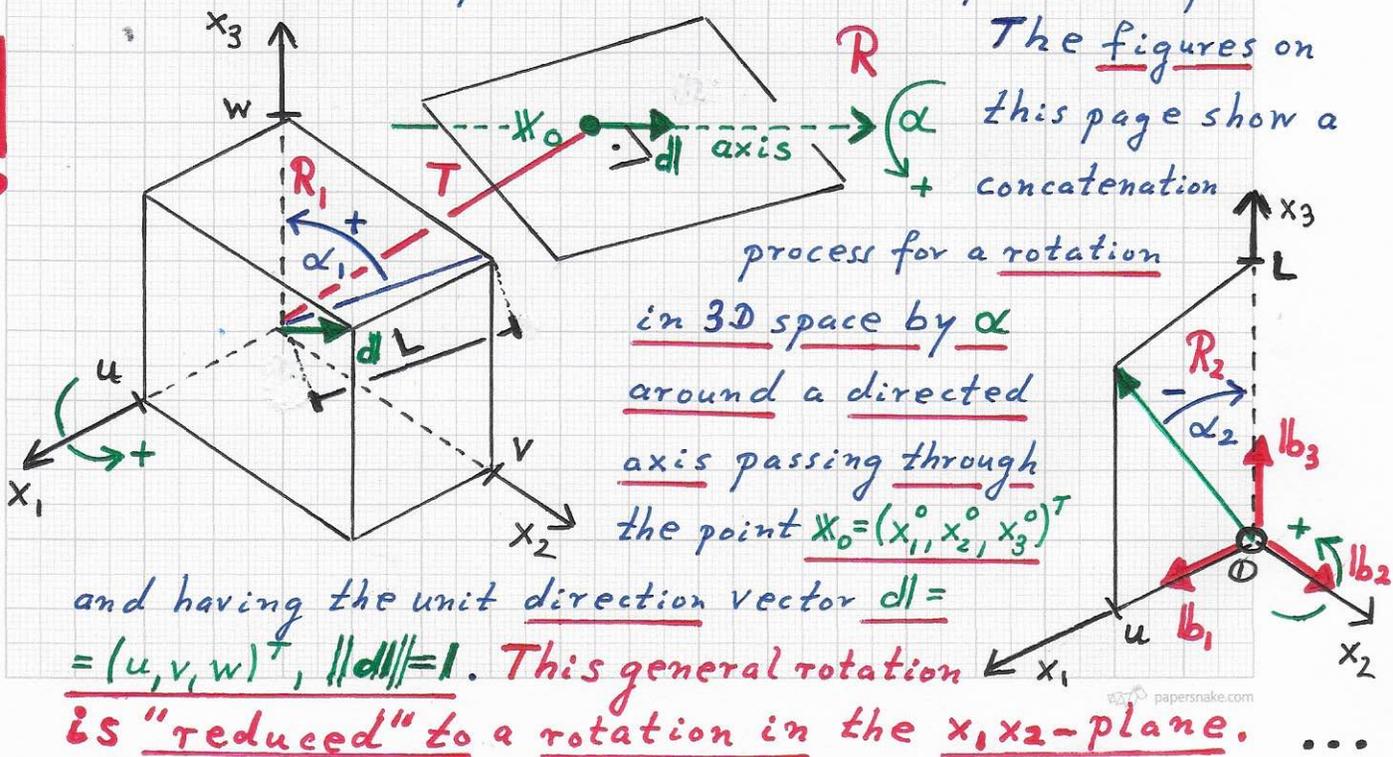


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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

• Laplacian eigenfunctions and neural networks:... We briefly review the "simplification" of an arbitrary rotation in 3D space by decomposing the desired rotation operation into "basic transformations" that are concatenated — as commonly done in computer graphics. The main underlying fact used in this context is the fact that a rotation is always an operation "taking place in a plane, i.e., a 2D subspace," regardless of the dimension of the n -dimensional embedding space ($n \geq 2$). Understanding a general rotation as a concatenation of "basic transformations" is a more intuitive approach — and the final, complicated matrix product must be computed only once.



OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

• Laplacian eigenfunctions and neural networks:...

The general rotation R is decomposed into the following sequence of "basic transformations":

- i) T : translation by positional vector $-x_0$
- ii) R_1 : rotation by α_1 around x_1 -axis (+rotation)
- iii) R_2 : rotation by $-\alpha_2$ around x_2 -axis (-rotation)
- iv) R_3 : rotation by α around x_3 -axis (+rotation)
- v) R_2^{-1} : inverse of rotation R_2
- vi) R_1^{-1} : inverse of rotation R_1
- vii) T^{-1} : inverse of translation T

If we view the names of these seven transformations as matrices, then the final general rotation matrix R will be $T^{-1}R_1^{-1}R_2^{-1}R_3R_2R_1T$.

It is assumed that all matrices are written in homogeneous form, i.e., as 4-by-4 matrices:

$$\underline{T} = \begin{bmatrix} 1 & 0 & 0 & -x_1^0 \\ 0 & 1 & 0 & -x_2^0 \\ 0 & 0 & 1 & -x_3^0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad \underline{T}^{-1} = \begin{bmatrix} 1 & 0 & 0 & x_1^0 \\ 0 & 1 & 0 & x_2^0 \\ 0 & 0 & 1 & x_3^0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad \underline{R}_1 = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & c\alpha_1 & -s\alpha_1 & 0 \\ 0 & s\alpha_1 & c\alpha_1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad \underline{R}_1^{-1} = R_1^T,$$

$$\underline{R}_2 = \begin{bmatrix} c(-\alpha_2) & 0 & +s(-\alpha_2) & 0 \\ 0 & 1 & 0 & 0 \\ -s(-\alpha_2) & 0 & c(-\alpha_2) & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} c\alpha_2 & 0 & -s\alpha_2 & 0 \\ 0 & 1 & 0 & 0 \\ s\alpha_2 & 0 & c\alpha_2 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad \underline{R}_2^{-1} = R_2^T,$$

$$\underline{R}_3 = \begin{bmatrix} c\alpha & -s\alpha & 0 & 0 \\ s\alpha & c\alpha & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}.$$

The figures on the previous page provide the needed values for α_1 and α_2 :

$$\underline{c\alpha_1} = \underline{\cos\alpha_1} = \underline{w/L}, \quad \underline{s\alpha_1} = \underline{\sin\alpha_1} = \underline{v/L};$$

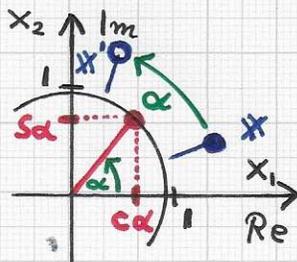
$$\underline{c\alpha_2} = \underline{\cos\alpha_2} = \underline{L/l} = \underline{L}, \quad \underline{s\alpha_2} = \underline{\sin\alpha_2} = \underline{u/l} = \underline{u}; \quad \text{where}$$

$$\underline{L} = \underline{(v^2 + w^2)^{1/2}}.$$

Stratoran■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

• Laplacian eigenfunctions and neural networks:...

The orthonormal basis vectors defining the 3D coordinate system (together with its origin \odot) are $\mathbf{l}_1, \mathbf{l}_2$ and \mathbf{l}_3 . The described concatenation method is based on the idea of using the plane spanned by vectors \mathbf{l}_1 and \mathbf{l}_2 as the plane (= 2D subspace) for executing the rotation, using the x_3 -axis with direction \mathbf{l}_3 as rotation axis. There exists a close relationship between this idea and using geometric algebra for performing general rotations in an n -dimensional space.



• Another relationship exists between complex numbers and geometric algebra in the context of rotations. The left figure shows how complex numbers can be used to describe rotations in the (complex) plane. The desired rotation "operator" is the complex number $x_R = c\alpha + i s\alpha$; the point/number to be rotated is $x = x_1 + i x_2$ ($\hat{=} (x_1, x_2)^T$); the image of x is $x' = x_R x = (c\alpha + i s\alpha)(x_1 + i x_2) = (x_1 c\alpha - x_2 s\alpha) + i(x_1 s\alpha + x_2 c\alpha)$. This result is equivalent to the result $\begin{pmatrix} c\alpha & -s\alpha \\ s\alpha & c\alpha \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} x_1 c\alpha - x_2 s\alpha \\ x_1 s\alpha + x_2 c\alpha \end{pmatrix}$.

In the following, these relationships will become clearer.

We will discuss the wedge product, multiplication rules and rotors. ...

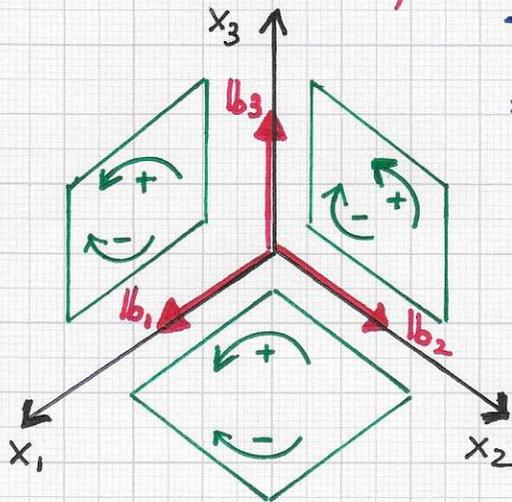
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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

• Laplacian eigenfunctions and neural networks:...

• We will now focus on geometric algebra and the wedge product,

to see the close relationships with the described related topics. The goal is an "intuitive understanding" of a ROTOR, a rotation operation that directly applies to n-dimensional space.



The left figure and multiplication table summarize the relationships between basis vectors lb_1, lb_2, lb_3 and basis bivectors $lb_{12}, lb_{23}, lb_{31}$; between wedge products and the associated parallelograms; and between orientation and sign of wedge products. Anti-commutativity yields the wedge products

\wedge	lb_1	lb_2	lb_3
lb_1	0	lb_{12}	$-lb_{31}$
lb_2	$-lb_{12}$	0	lb_{23}
lb_3	lb_{31}	$-lb_{23}$	0

$$lb_{12} = lb_1 \wedge lb_2 = -lb_{21}$$

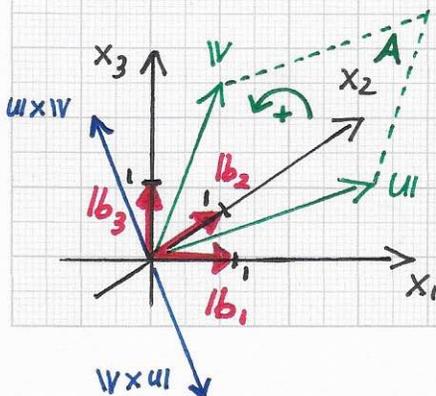
$$lb_{23} = lb_2 \wedge lb_3 = -lb_{32}$$

$$lb_{31} = lb_3 \wedge lb_1 = -lb_{13}$$

$$(u_i \wedge u_j) = -(u_j \wedge u_i)$$

These rules can now be used to

compute the wedge product of two vectors $u_i = (u_1, u_2, u_3)^T$ and $v_i = (v_1, v_2, v_3)^T$:



$$\begin{aligned} u_i \wedge v_i &= (u_1 lb_1 + u_2 lb_2 + u_3 lb_3) \wedge (v_1 lb_1 + v_2 lb_2 + v_3 lb_3) = \\ &= u_1 v_1 lb_1^2 + u_1 v_2 lb_1 lb_2 + u_1 v_3 lb_1 lb_3 + u_2 v_1 lb_2 lb_1 + u_2 v_2 lb_2^2 + u_2 v_3 lb_2 lb_3 \\ &+ u_3 v_1 lb_3 lb_1 + u_3 v_2 lb_3 lb_2 + u_3 v_3 lb_3^2 = (u_1 v_2 - u_2 v_1) lb_{12} + \dots \\ &+ (-u_1 v_3 + u_3 v_1) lb_{31} + (u_2 v_3 - u_3 v_2) lb_{23} \dots \end{aligned}$$

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■ OBJECT AND MATERIAL EIGENFUNCTIONS - Cont'd.

• Laplacian eigenfunctions and neural networks:...

The figure at the bottom of the previous page also includes the

cross products $u_i \times v_j$ and $v_j \times u_i$. The cross product $u_i \times v_j$ is

$$\begin{vmatrix} u_i & v_j & k \\ u_2 & v_2 & j \\ u_3 & v_3 & i \end{vmatrix} = +i(u_2 v_3 - u_3 v_2) + j(-u_1 v_3 + u_3 v_1) + k(u_1 v_2 - u_2 v_1) \hat{=} \begin{bmatrix} u_2 v_3 - u_3 v_2 \\ -u_1 v_3 + u_3 v_1 \\ u_1 v_2 - u_2 v_1 \end{bmatrix} = u_i \times v_j.$$

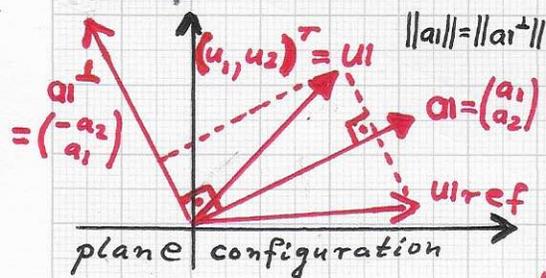
Thus, $\|u_i \wedge v_j\| = \|u_i \times v_j\| = \text{area of parallelogram}(A)$.

- We point out again that $u_i v_j = \langle u_i, v_j \rangle + (u_i \wedge v_j)$, implying that the orthonormal basis vectors satisfy $b_i^2 = \langle b_i, b_i \rangle = 1$ and $b_i b_j = 0 + (b_i \wedge b_j) = (b_i \wedge b_j) = b_{ij} = -b_{ji}$, $i \neq j$.

\wedge	b_1	b_2	b_3	b_{12}	b_{23}	b_{31}	b_{123}
b_1	1	b_{12}	$-b_{31}$	b_2	b_{123}	$-b_3$	b_{23}
b_2	$-b_{12}$	1	b_{23}	$-b_1$	b_3	b_{123}	b_{31}
b_3	b_{31}	$-b_{23}$	1	b_{123}	$-b_2$	b_1	b_{12}
b_{12}	$-b_2$	b_1	b_{123}	-1	$-b_{31}$	b_{23}	$-b_3$
b_{23}	b_{123}	$-b_3$	b_2	b_{31}	-1	$-b_{12}$	$-b_1$
b_{31}	b_3	b_{123}	$-b_1$	$-b_{23}$	b_{12}	-1	$-b_2$
b_{123}	b_{23}	b_{31}	b_{12}	$-b_3$	$-b_1$	$-b_2$	-1
b_4	$-b_{14}$	$-b_{24}$
b_5	$-b_{15}$	$-b_{25}$

The left table lists

the possible wedge products for $b_1, b_2, \dots, b_3, b_{123}$ (and indicates how to extend the table). The



Left figure shows the elementary operation of a ROTOR, the "perpendicular reflection" of u_i relative to axis vector a_i :

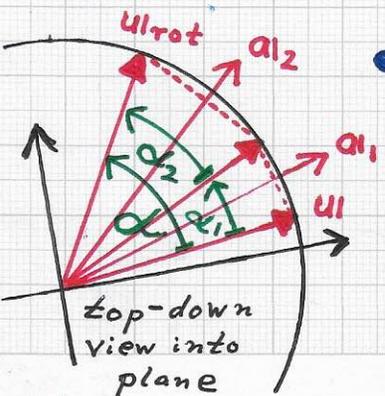
$$u_{i,ref} = u_i - 2 \frac{\langle u_i, a_i \rangle}{\|a_i\|^2} \cdot a_i = \dots = a_i u_i a_i^{-1}.$$

- As discussed earlier, we can achieve a rotation by an angle via two reflections, see left figure.

IN GEOMETRIC ALGEBRA, THIS ROTATION IS EXPRESSED AS

$$u_{i,rot} = a_{i2} a_{i1} u_i a_{i1}^{-1} a_{i2}^{-1} = R u_i R^{-1}.$$

THE ROTOR IS CALLED R.



set $\alpha_1 = \alpha_2 \Rightarrow \alpha = 2\alpha_1$

$$R = a_{i2} a_{i1}$$

$$= \|a_{i2}\| \|a_{i1}\| \cos \frac{\alpha}{2} + \|a_{i2}\| \|a_{i1}\| \sin \frac{\alpha}{2} \cdot \frac{(a_{i2} \wedge a_{i1})}{\|a_{i2} \wedge a_{i1}\|}$$